Einstein's Theory of Gravity and the Problem of Missing Mass

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The observed matter in the universe accounts for just 5% of the observed gravity. A possible explanation is that Newton's and Einstein's theories of gravity fail where gravity is either weak or enhanced. The modi ed theory of Newtonian dynamics (MOND) reproduces, without dark matter, spiral-galaxy orbital motions and the relation between luminosity and rotation in galaxies, although not in clusters. Recent extensions of Einstein's theory are theoretically more complete. They inevitably include dark elds that seed structure growth, and they may explain recent weak lensing data. However, the presence of dark elds reduces calculability and comes at the expense of the original MOND premise { that the matter we see is the sole source of gravity. Observational tests of the relic radiation, weak lensing, and the growth of structure may distinguish modi ed gravity from dark matter.

Introduction

The problem of missing mass has been with us more than 70 years: Given the amount of directly observable matter, general relativity (GR, Einstein's theory of gravity) produces too little gravity to account for a host of observations. On scales of one to tens of kiloparsecs, the observed random or coherent velocities of stars and gas are much greater than the escape velocity in the self-gravity of those same stars, gas and dust. The same is true for galaxy clusters on much larger scales. Gravitational potentials around galaxy clusters, deeper than would be produced by the observed matter, are also needed to explain observed gravitational lensing; that is, the deformation of light bundles from background galaxies.

Evidence also exists for anom alously strong gravity on the largest observable scale: out to the cosm ological horizon. In a universe that contained only ordinary matter (often called baryonic matter, encompassing protons and neutrons, which make up over 99.9% of the mass of ordinary matter), the growth of structures, such as galaxies and clusters of galaxies, would be suppressed. During recombination, when that universe was approximately 400;000 years old, the seeds for galaxies and clusters would be entirely erased by dissipation (known as Silk damping) and no structure would form on scales up to many tens of megaparsecs.

The now standard solution of these dynamical mysteries is that an unobserved form of mass, which exceeds the observed mass of both galaxies and clusters,

provides the gravity that prevents them from ying apart, increases lensing and prevents Silk damping. The missing mass neither shines nor absorbs or scatters light enough to be directly detected by our telescopes. It should be close to pressureless and must be non relativistic well before recombination. It is therefore known as cold dark matter (CDM). A number of candidate particles have been proposed that have these properties and their behaviour in the Universe has been studied in exquisite detail.

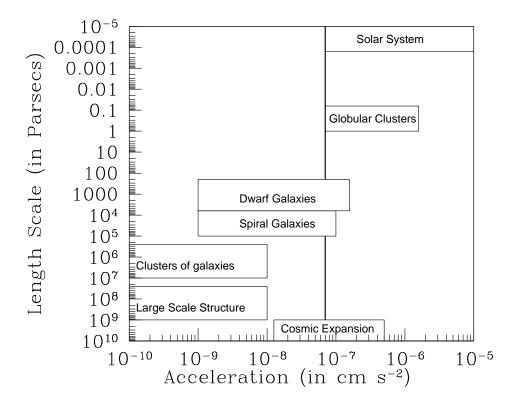


Figure 1: Evidence for dark matter or for deviations from GR tend to appear in systems in which the acceleration scale is weak (to the left of the solid horizontal line) at about 7 $\,$ 10 8 cm $\,$ 2 . There is strong evidence for either of the above in dwarf galaxies, spiral galaxies, clusters of galaxies, the large scale structure of the Universe and in the expansion of the Universe itself.

An altogether di erent approach can be taken if one notes that the evidence for m issing m ass arises because of a m ism atch between the gravitational eld one would predict from the observed m ass distribution in the Universe and the observed gravitational eld. The observed discrepancies arise when the elective gravitational acceleration is around, or below, a_0 ' 10 8 cm sec 2; that is in a regime of very weak gravitational eld. Perhaps the Newtonian theory of gravity—and GR-break down in this regime. In this Review we provide an updated assessment of this theory.

M odifying Newtonian Gravity

The possibility that Newtonian gravity and GR do not accurately describe very weak gravitational elds was proposed more than 25 years ago. Milgrom suggested that Newton's second law, F = m a (where F is the gravitational force applied to a unit of mass m to produce an acceleration a) is modiled when gravity is weak, to F = m ($a_j = a_0$) a [1]. This proposal has been named modiled Newtonian dynamics (MOND). More modern versions of MOND cast it instead as a modiled theory of gravity, altering the Newton-Poisson equation that relates the gravitational force to the distribution of mass density that is responsible for it [2].

MOND has a number of appealing features. It independently explains the empirical Tully-F isher relation between the luminosity, L, of a spiral galaxy and its asymptotic rotational velocity, $v: L / v^4$ where G is the gravitational constant and M is the baryonic mass of the galaxy. G iven the ratio of the baryonic mass of these spirals to their luminosity (the mass-to-light ratio), this is equivalent to $(a_0G)M = v^4$, exactly what would be predicted by MOND. A systematic study of a wide range of spiral galaxies pins the acceleration scale to be unique: a_0 ' 12 10 8 cm s ². Furthermore, the detailed features of the rotational velocity as a function of radius are predicted by the baryonic mass distribution [3]. MOND has also been used to predict the analogue of the Tully-F isher relation for elliptical galaxies (the Faber-Jackson relation between baryonic mass and velocity dispersion), the existence of galaxies with low surface brightness, and an upper limit on the mean surface brightness of spiral galaxies (known as Freeman's law) [4].

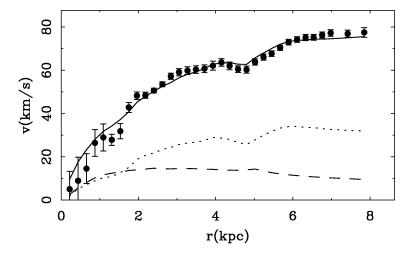


Figure 2: The MOND rotation curve of the galaxy NGC1650 (solid line) reproduces observed features with surprising delity given just one free parameter { the mass-to-light ratio M = L = 0.43. Also shown is the Newtonian rotation curve that would result from just the gas (dotted) or just the stars (dashed). The quality of the rotation curve the from MOND is generic. (Figure courtesy of M caugh; originally published in [42])

Attempts to resolve the mass discrepancy on the scale of clusters of galaxies

have been more problem atic. One is obliged either to use a value of a_0 that is different from the one used for galaxies or to assume the existence of a small amount of dark matter. If one of the three types of neutrino (electron, muon or tau) has a mass of a few electrons volts it would be an ideal candidate for cluster dark matter in MOND [5].

An added complication is that gravitating systems cannot be studied in isolation and the external gravitational eld can play a role in the internal dynamics of disparate objects such as star clusters, molecular clouds and galaxies. MOND does not satisfy Birkho 's theorem (the analogue in gravity of Gauss' law in electrom agnetism) for real masses [6, 7, 8] and this means that the acceleration of any real probe { a star or a cloud of gas { even if it is located in a spherically symmetric system, depends not just on the mass that is interior to the probe but on the mass that is exterior as well.

MOND was developed as a phenom enological description of spherically (or at least axially) sym metric, non-relativistic, low acceleration systems. Until a fully dynamical relativistic theory can be constructed, MOND itself cannot reliably predict, among other things: anomalous accelerations in galaxy clusters, the ects of gravitational lensing of light, the expansion of the Universe or the growth of structure.

Relativistic theories of modied gravity

Despite MOND's successes and failures, for it to be seriously judged as a candidate explanation of anything, it must be embedded in a modication of Einstein's GR theory. Einstein cast gravity as a geometric theory of spacetime (the combination of space and time). The properties of spacetime are encoded in a 4 by 4 symmetric matrix with 10 free components, which is called the metric, g, which is itself a function of space and time. From this metric, one can construct various geometric quantities, such as the overall curvature of spacetime, R, known as the Ricci scalar, the Ricci tensor R and the Riemann tensor, R.

E instein proposed that the energy content of the Universe would source these various quantities, curving spacetime according to a xed set of rules, called the Einstein eld equations. The dierent components of the Universe would in turn respond to the curvature of spacetime: In the absence of non gravitational forces, they would follow geodesic paths that could be derived from the metric.

E instein's theory must be tampered with to incorporate MOND. There are two possible ways ofmodifying it. One way is to change how curvature responds to the presence of matter. The rules that E instein posited for deriving gravity start with the simplest of all the quantities that encode curvature, the Ricci scalar. A

rst step away from E instein is to bring in other functions of the metric such as the Ricci and Riemann tensors, as well as more complicated functions of the Ricci scalar. Indeed, because of the elects of quantum mechanics on the spacetime metric, one expects the full theory of gravity to be a more complex combination of various geometric quantities [9]. As a result, one inevitably introduces new gravitational degrees of freedom.

Some of these modications have consequences at short distance scales and result in small (although potentially measurable) corrections to standard physics, which are insuccient to reproduce galactic rotation curves. Others result in modications to cosmology, but not in MOND-like behavior on galactic scales. Nevertheless a few proposals have been advanced for modications of gravity that can play a role on galactic and supra-galactic scales. In [10] the Ricci tensor is replaced by the Weyl curvature and a scalar eld is introduced to play the role of a variable Newton's constant. In [11], a logarithm of the Ricci curvature is considered as the fundamental action. For these cases and others, there is an extensive program to explore the theoretical and phenomenological consequences.

A complementary approach is to postulate that light and matter respond to the geometry of space and time dierently than predicted by Einstein. The simplest way to implement this is to distinguish the geometric metric which responds dynamically to the contents of the Universe from the geodesic metric which dictates (in the absence of other, non-gravitational, forces) how those contents propagate through spacetime. The simplest such theory relates the two metrics by a location-dependent change of scale, known as a conformal transformation, and endows the scalar eld describing this transformation with a dynamics of its own. This theory has been extensively studied in many contexts and is highly constrained [12]. Whereas such a scalar eld can a ect the dynamics of massive bodies, it doesn't modify the propagation of light rays and therefore will not play a role in phenomena such as gravitational lensing.

More general transform ations involve introducing not only a change in scale between the metrics, but also a distortion of angles and this can be done, for example, by introducing a preferred time direction—or a preferred rest frame. The most natural implementation is to add in a spacetime vector—eld that has a non-zero value at each point in spacetime; in other words to point in some direction in spacetime. If that direction is chosen to be (on average) the direction that denes the future (forward in time) as opposed to some direction in space, then the preferred direction will be established [13].

Bekenstein [14] recently proposed a fully relativistic theory that included all of these elements: a disformal relation between the geometric and geodesic metrics, a preferred frame, and modied dynamics for the geometric metric. For an appropriate choice of an arbitrary but universal function, his theory could lead to MOND ian dynamics on galactic scales. Bekenstein's theory is known as TeVeS, where the T stands for tensor (representing the metric), V for the time-like vector eld, and S for the scalar eld responsible for the scale transformation. TeVeS is actually part of a wider class of models that reproduce MOND on astrophysical scales. An alternative subclass of these models dubbed generalized Einstein-Aether (GEA) theories, include only the time-like vector eld and no scalar eld [15].

These modi cations of Einstein GR are su ciently well de ned that it is possible to make m predictions within each model for what should be observed on various astrophysical and cosmological scales. Such theories have interesting properties, but whereas Einstein's original proposal is highly constrained, these more complex proposals are less so. Inevitably, they involve extra elds that may

Observational tests and limitations

With a relativistic theory of modi ed gravity in hand, it is possible to make a number of predictions on a range of scales. For a start, one can focus on the e ect that TeVeS or GEA will have on the gravitational eld of compact objects, such as stars or black holes. It has been shown that the atom ic spectral lines from the surface of stars will be a ected by TeVeS parameters [6] whereas farther out, there may be directly detectable preferred frame e ects that will modify the Newtonian orbits of nearby objects [17]. On even larger scales, it has been proposed that the dierence in ight time between gravity waves and neutrinos from, for example, a supernova can be used as a signature of modi ed gravity theories [18]. As yet, an analysis of millisecond binary pulsars, one of the GR laboratories par excellence, is still lacking.

Relativistic theories of modi ed gravity make speci c predictions about the dynam ics of the Universe. The TeVeS theory has a particular property: The energy density in the extra elds is always proportional to the energy density of whatever is the dominant form. Furthermore, the constant of proportionality is independent of the initial conditions and dependent solely on the fundamental constants of the theory [19]. These features lead to a tight constraint on the overall energy density in the extra elds { it cannot be more than one fith of the contribution from baryons. Thus, unlike dark matter, this energy density is subdominant to the baryonic mass and does not a ect the overall expansion rate. Such behavior can be found in other proposals for modi ed gravity but it is not generic. It is challenging to not a parametrization replacing CDM [20], but such theories can predict a wide range of cosmological behavior, from the highly regular to the unstable, leading to accelerated expansion or to contraction on a nite time scale [21, 22].

M uch of the recent advance in cosm ology has been acom plished through understanding and measuring the statistical properties of the growth and morphology of large scale structure, through the cosmic microwave background (CMB) and through surveys of galaxies. With relativistic theories of modied gravity it is now possible to make predictions on the largest scales, and this has been done for a selection of the currently proposed models, in particular for the original TeVeS model and for GEA theories.

Inhom ogeneities evolve in a more complex way in these theories than in the case of GR, with two main new qualitative features. First, the extra degrees of freedom drive the initial growth of perturbations; they seem to be a necessary piece of the theory and there seems to be no other way to seed structure given the constraints from observations of the CMB on the amplitude of uctuations in the baryonic matter density when the universe was 1000 times smaller than it is today. Gravity alone, even (stronger) MOND ian gravity, appears to be incapable of growing structure without seeds of such structure that are less coupled to the photons than are baryons [6]. TeVeS and GEA avert this conundrum by

allowing modes of the new gravitational elds to grow and seed baryonic structure. E ectively, these new degrees of freedom act as dark elds 19, 23].

A prelim inary comparison between the TeVeS theory and both the CMB and large scale structure data indicates that they are roughly compatible [19]. There are a few caveats. First, it may be necessary to include a non-negligible amount of massive neutrinos with a mass of a few electron volts. This is also the mass range required by MOND to agree with clusters. It is still unclear whether this is generic [24], but if indeed it is, it may be testable in the near future. Experiments such as KATRIN (the Karlsruhe tritium neutrino experiment) will bring constraints on the mass of the neutrinos to below 1 electron volt [25].

Second, there is a subtle e ect that can emerge on the largest scales. In GR, when most of the matter is non-relativistic (in the form of atoms or dark m atter), perturbations in the metric can be described in terms of one function, which on small scales is the Newtonian potential that gives rise to the inverse square law of gravity. In modi ed theories of gravity, perturbations in the metric are generally described in term s of two potentials, one of which is the Newtonian potential (apart from the deviations required to lead to MOND). The di erence between the two potentials, also known as gravitational slip, can lead to changes in the growth of structure, large-scale gravitational lensing and anisotropies in the CMB. For example, it is still unclear whether it is possible to completely m atch both the CMB data on large and small scales as well as the amplitude of mass uctuations in galaxy surveys. If one is to boost the small angle peaks of the predicted angular power spectrum of the CMB so that it can match the data, one runs the risk of introducing large uctuations on large scales due to the gravitational slip. Furtherm ore this e ect can suppress the amount of clustering on galactic, cluster and super cluster scales. So as yet, the comparison between TeVeS, the CMB and large scale structure is not conclusive [26], and all the more so for GEA and other modications of gravity.

The gravitational slip may be the smoking qun for modied theories of gravity. There are, by now, a few suggestions on how it may be detected. The idea is simple: Die erent cosmological data probe dierent combinations of the two gravitational potentials and by combining them, it may be possible to tease out evidence for the slip. So, for example, a galaxy catalogue will be a measurement of the density contrast of the Universe and will be directly related to one of the potentials, whereas a map of large scale ows should probe the other potential. Measurements of weak lensing will depend on the sum of the two potentials as will observations of the CMB. By cross-correlating maps of the CMB with galaxy surveys, or alternatively matching maps of weak lensing with peculiar velocity ows it should be possible to search for gravitational slip, and if found, it will give us information on about the most relevant modi cations to theories of gravity [6, 27, 28]. Relativistic theories of modi ed gravity can be used to calculate the e ects of gravitational lensing and can be tested with the many well-measured gravitational lenses. A notable example is the Bullet cluster where the baryonic m ass, which is usually relatively well traced by hot x-ray em itting gas, is severely m isaligned with the sources of gravitational lensing as inferred from the distorted im ages of background galaxies [29]. A dark matter explanation for the Bullet

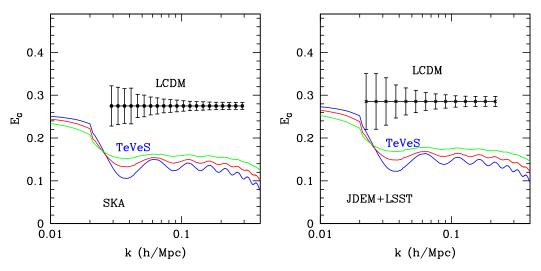


Figure 3: The normalized cross correlation spectrum, $E_{\rm G}$, between density and weak lensing on large scales, as a function of wave number, k. The points and errors bars are a forecast of what would be expected in a Universe with dark matter, a cosmological constant and Einstein gravity (LCDM) as measured by the Square Kilometre Array (SKA) or by a combination of the Large Synoptic Survey Telescope (LSST) and a possible version of the Joint Dark Energy Mission (JDEM) [27]. The coloured lines are variants of the TeVeS model of modiled gravity. The two dilerent classes of theories are clearly distinguishable (Figure courtesy of P. Zhang, drawn from [27]).

cluster is that the lensing is centered on two localized accumulations of dark matter, the halos of two colliding clusters that have passed through each other, and the hot cluster gas is interacting as a result of that collision.

It would appear dicult to reconstruct such a con guration merely by modifying gravity, but two features of such theories prevent such a simple assessment. The rst is the presence of the extra dynamical elds. Because structure is seeded in these models by uctuations in these extra elds, unsourced by any baryonic uctuations, these elds clearly are capable of supporting uctuations that are a source of gravity independent of the baryons. In principle, these gravity—eld seeds may evolve into dark—eld concentrations that interact only gravitation—ally, become separated from their associated baryons in a collision, and source gravitational lensing as seen in the bullet cluster.

The second feature is that these theories do not satisfy Birkho 's theorem. As a result, not only is the gravitational eld due to a localized concentration of matter that is not unique (and may depend on the history that led to its assembly) but the environment can play a major role in the interactions between two systems. For example, in inferring the dynamics of galaxies within a cluster it becomes necessary to include the elects from the rest of the cluster, from neighbouring clusters and from any enveloping super-cluster [8, 7]. Yet, these theories should be predictive and there have been attempts to study lensing properties of special capacities and galaxy clusters in TeVeS and GEA, with mixed results [30, 31]. Without including the elect of extra degrees of freedom or the environ-

ment, it was found that, in TeVeS, the Bullet cluster would have to be surrounded by massive neutrinos, with a mass of approximately 2 electron volts [32]. Even with the inclusion of the electron of the extra degrees of freedom, it would be necessary to have some form of dark matter, which could be in the form of neutrinos. This is all hardly surprising because, as discussed above, MOND requires clusters to have some dark matter. For GEA theories it is possible in principle to reconstruct the geometry and gravitational eld of a lens such as the Bullet cluster, without any extra matter but with a substantial contribution from the extra degrees of freedom [33]. The question of whether the dynamical evolution of the dark eld perturbations leads naturally to such large dark eld halos at large time excales remains outstanding.

A Ithough there is nothing intrinsically inconsistent with having the new elds that mediate the modications of gravity envisioned in MOND act as dark seeds of structure or dark concentrations of gravitational lensing, this necessity detracts from the cleanliness of the original MOND vision: What you see is apparently not what you get, even in MOND. It is therefore much harder to make testable predictions for modied theories of gravity than was already thought and thus far, these relativistic extensions of MOND remain viable solutions to the problem of missing mass.

Dark Energy and future measurements

The problem of m issing gravity has been at the forefront of cosm ology for m any decades. From the moment a proposal was put forward to solve the missing mass problem of galaxies with modiled gravity, it was realized that there could be cosm ological implications: The acceleration scale, a_0 ' $10^8 \, \text{cm} = \text{s}^2$, which characterizes the transition between Newtonian and non-Newtonian gravity in MOND, is of the same order of magnitude as cH_0 , where c is the speed of light and H $_0$ is the expansion rate of the Universe today.

Furtherm ore, the recent discovery that the Universe m ay be accelerating (and not decelerating as GR predicts if the energy content of the Universe is dominated by pressureless, non-relativistic matter) has been taken to imply the existence of an additional dark component of the Universe. Dubbed dark energy it is gravitationally repulsive and can drive the expansion at late times. It may be possible that the accelerated expansion of the Universe is instead due to modications to E instein's theory of gravity.

Some modi ed gravity theories seek to explain only the accelerated expansion and not the missing mass problem of galaxies and clusters. DGP (D vali-G abadadze-Porrati) [34] and (generic) f (R) theories [35] are two that have received widespread attention. However, as for united models of the dark sector, where an extra component of the energy density of the Universe can be both dark matter and dark energy, it is natural to consider models where modications of gravity can give rise to both the breakdown of Newtonian gravity on galactic scales and to the acceleration of cosmic expansion, as in TeVeS or GEA [21, 37, 15, 36, 38].

As argued in the previous section, future surveys may be able to distinguish between theories of modi ed gravity and GR with dark matter and dark energy [27]. A survey of galaxies (such as those that will be done by the Joint Dark Energy Mission [39] or the Square Kilometre Array [40]) combined with a measurement of lensing on large scales (such the one proposed by the Joint Dark Energy Mission or the Large Scale Synpotic Survey [41]), should be able to clearly pick out the signature of gravitational slip of a theory such as TeVeS. Furthermore, this would be on scales for which many of the issues that complicate predictions in the case of quasi-isolated systems such as clusters, would not come into play.

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