

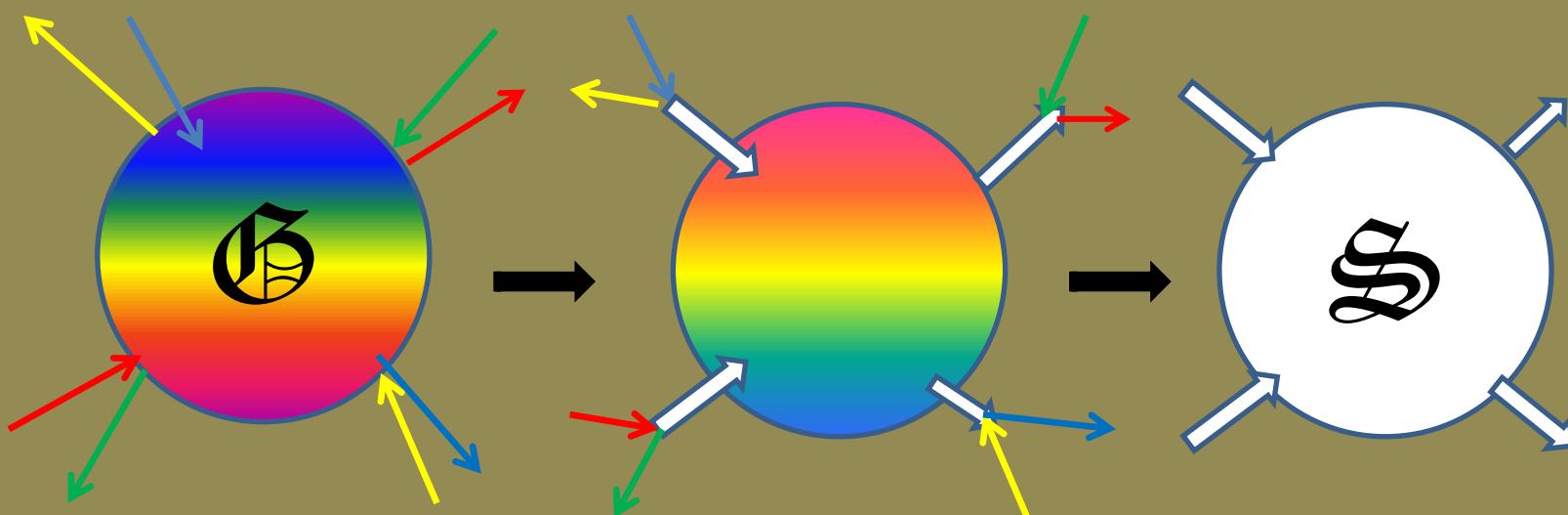
# **ЯСИЯ ПО ТОТИЧЕСКИЕ РЕЖИМЫ АДРОННЫХЪ ПРОЦЕССОВЪ ВЪ КУД**

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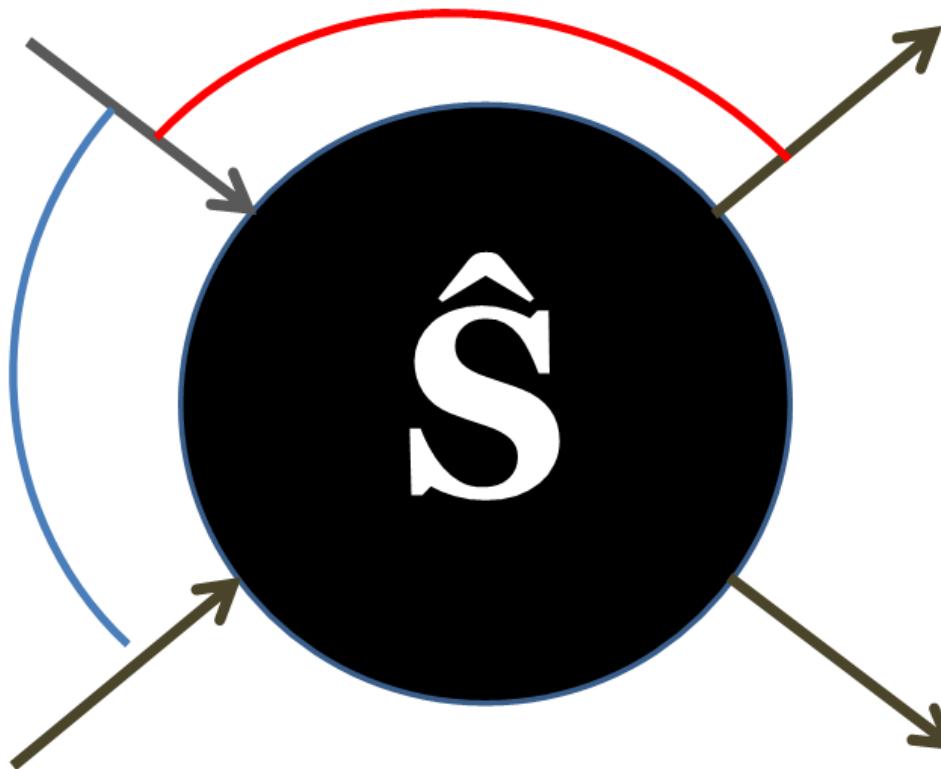
# ОТЪ КВАРКЪ - ГЛЮОННЫХЪ

## ФУНКЦИИ ГРИНА

### КЪ АДРОННОИ S-МАТРИЦѢ



# Scattering Regimes



Hard:

$$s \gg m^2; \frac{s}{t} \text{ fixed}$$

“Short distances”

Semi-Hard

$$m^2 \ll -t \ll s$$

Soft:

$$s \gg m^2; \quad 0 \leq -t \ll s$$

“Large distances”

# Amplitude

Massless QCD

$$T(s + i0, t) = F(s, t; \Lambda^2, \{c_h\})$$

$$c_h = \frac{M_h^2}{\Lambda^2}$$

$$\Lambda^2 = \mu^2 \exp\left(- \int^{g^2} \frac{dx}{\beta(x)}\right)$$

$$\Lambda^2|_{g^2 \rightarrow 0} \sim e^{-1/(\beta_0 g^2)}$$

# Boundary Condition

$$\widehat{S}|_{g^2 \rightarrow 0} \rightarrow \widehat{1}$$

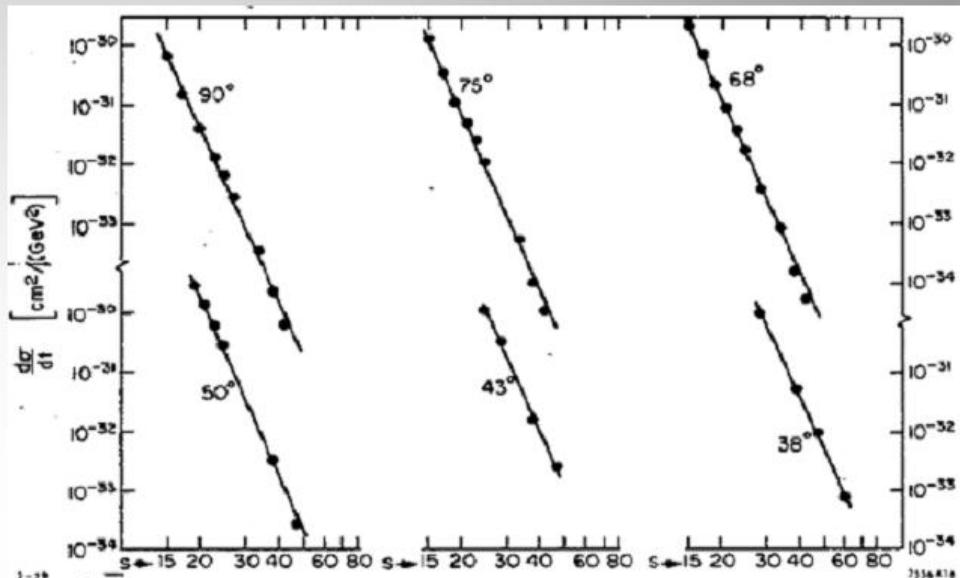
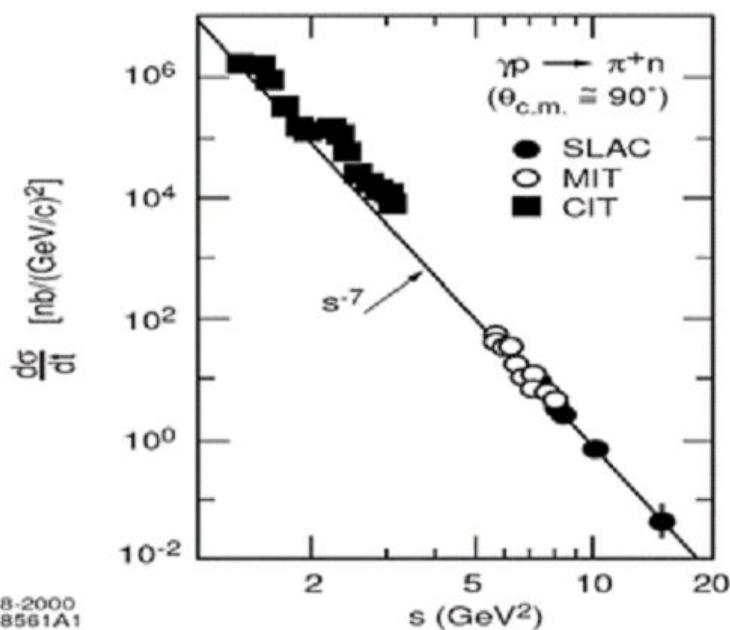
$$T|_{g^2 \rightarrow 0} \rightarrow 0$$

# Hard Scattering

$$\lim_{\substack{t \\ (s \rightarrow \infty, \frac{t}{s} \text{ fix})}} T(s, t) = \lim_{\substack{t \\ (s \rightarrow \infty, \frac{t}{s} \text{ fix})}} F\left(\frac{s}{\Lambda^2}, \frac{t}{s}\right) = F\left(\infty, \frac{t}{s}\right)$$

$$\lim_{(g^2 \rightarrow 0)} T(s, t) = \lim_{(\Lambda^2 \rightarrow 0, \frac{t}{s} \text{ fix})} F\left(\frac{s}{\Lambda^2}, \frac{t}{s}\right) = F\left(\infty, \frac{t}{s}\right) = 0$$

# Hard Scattering: Data



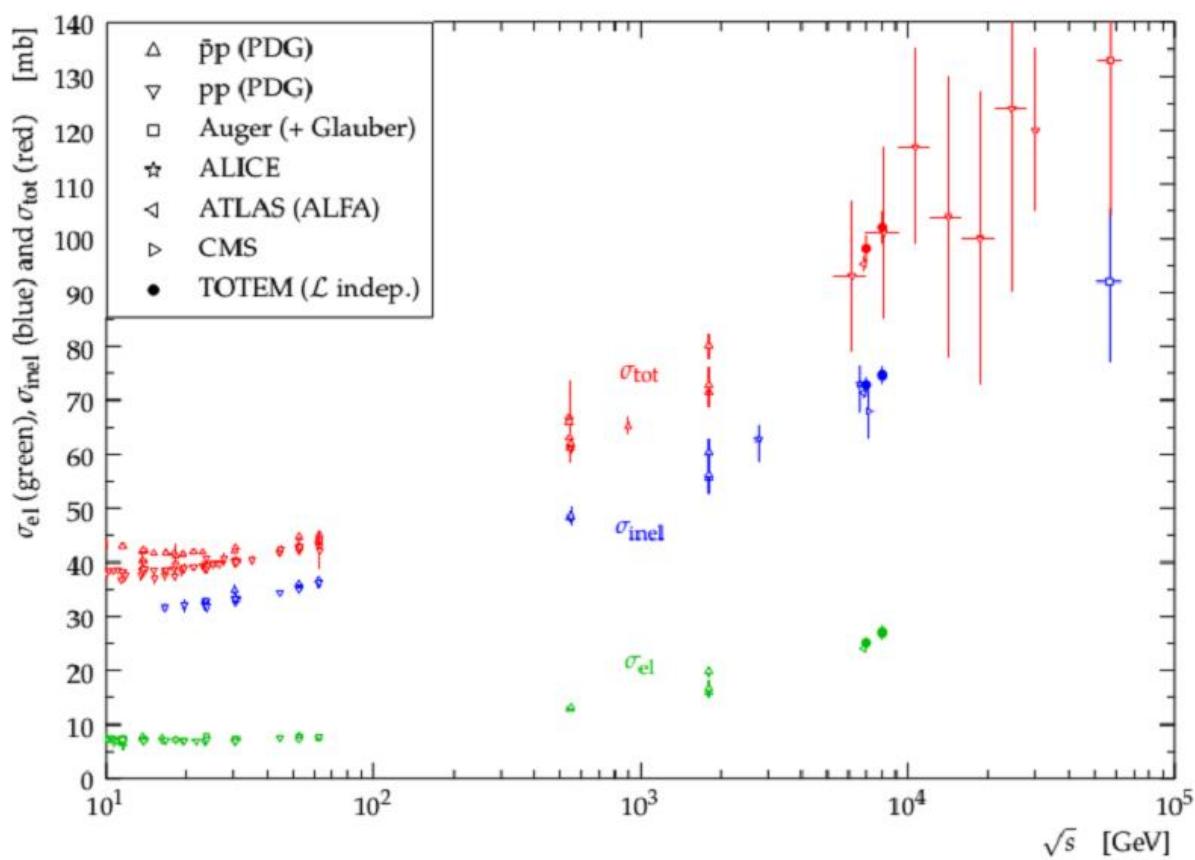
Constituent Counting Rule

Matveev V.A., Muradyan R.M., Tavkhelidze A.N. *Lett. Nuovo Cimento* 7, 719 (1973);

Brodsky S., Farrar G. *Phys. Rev. Lett.* 31, 1153 (1973)

$$T(s, t) \Big|_{(s \rightarrow \infty; \frac{t}{s} \text{ fix})} \approx \left( \frac{\Lambda^2}{s} \right)^N f\left(\frac{t}{s}\right)$$

# Soft Scattering: Data



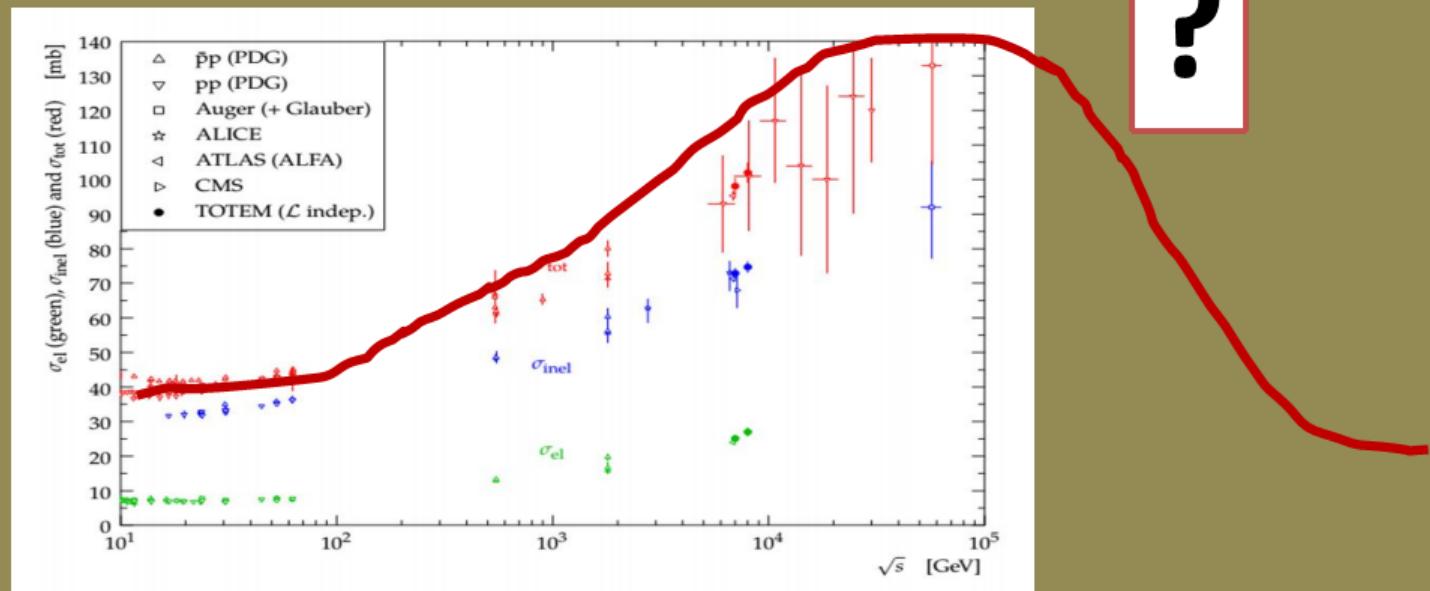
$$\sigma_{\text{tot}}(s) = \lim_{t \rightarrow 0} \frac{T(s + i0, t) - T(s - i0, t)}{2is}$$

# Soft Scattering

$$T(s, 0)|_{g^2 \rightarrow 0} \rightarrow 0$$

$$\lim_{g^2 \rightarrow 0} T(s, 0) = \lim_{\Lambda^2 \rightarrow 0} F(s/\Lambda^2) = \lim_{s \rightarrow \infty} F(s/\Lambda^2) = 0$$

$$\lim_{s \rightarrow \infty} \sigma_{tot}(s) = 0$$



## Adding mass

Let  $\exists i, i \in (1, 2, \dots, N_f) : m_{qi} \neq 0$

2nd mass scale

$$M^2 = m_i^2 \exp(L(g^2))$$

$$\frac{dL}{dg^2} = -\frac{\gamma_m}{\beta} = \frac{\gamma_{m0}g^2 + \dots}{\beta_0 g^4 + \dots}$$

$$\gamma_m = -\frac{\mu^2}{m^2} \frac{dm^2}{d\mu^2}$$

$$M^2 \sim \frac{1}{(g^2)^{\gamma_{m0}/\beta_0}}$$

# Quark Mass to Open the Way to Infinity ?

$$\lim_{g^2 \rightarrow 0} T(s, 0) = \lim_{\Lambda^2 \rightarrow 0} F\left(\frac{s}{\Lambda^2}, \frac{M^2}{\Lambda^2}\right) = F(\infty, \infty) = 0$$

$$\lim_{s \rightarrow \infty} T(s, 0) = \lim_{s \rightarrow \infty} F\left(\frac{s}{\Lambda^2}, \frac{M^2}{\Lambda^2}\right) = F\left(\infty, \frac{M^2}{\Lambda^2}\right)$$

A Toy Model

$$v > 1$$

$$Im T(s, 0) = \left(\frac{\Lambda^2}{M^2}\right)^v \frac{s}{\Lambda^2} \ln^2\left(\frac{s}{\Lambda^2}\right)$$

$$\sigma_{tot}(s) = \frac{1}{\Lambda^2} \left(\frac{\Lambda^2}{M^2}\right)^v \ln^2\left(\frac{s}{\Lambda^2}\right)$$

# Infrared fixed point

$\beta(g)$

Dilatation current:

$$D^\mu = T^{\mu\nu}x_\nu \quad \partial_\mu D^\mu = T$$

$$T = g_{\mu\nu}T^{\mu\nu} = \frac{\beta(g)}{g} F^2$$

0

$g$

$$g_* = \bar{g}(0)$$

$$\partial_\mu D^\mu|_{g=g_*} = 0$$

Scale invariance?

$$\Lambda^2|_{g \uparrow g_*} \approx (g_* - g)^{-1/\beta'_*} \quad \text{RG: } \mathfrak{D}(q^2)|_{q^2 \rightarrow 0} \sim \frac{1}{q^2} \left(\frac{\Lambda^2}{q^2}\right)^{\gamma_A(g_*)}$$

$$F(s, t)|_{g=g_*} = F(0, 0) = \text{const}$$

# ЗАКЛЮЧЕНИЕ

- ВЪ ЧИСТОЙ ГАУДИНАРИИ НЕОГРАНИЧЕННЫЙ РОСТЬ СВЕЧЕНИЙ ГАУБОНОВЪ НЕВОЗМОЖЕНЪ.
- ВЪ КХД СЪ ХОТА БЫ ОДНИМЪ ЖАССИЕНЫЯТЬ КВАРКОЖЪ, ИНДУЦИРОВАНИЯТЬ ВЗАИМОДЕЙСТВІЕ ІСІНЪ С ПОЯСЬЮ ХИГГСА, ТАКОЙ РОСТЬ НЕ ЗАПРЕЩЕНЪ.